

Using Bernoulli Equation to Solve Burger's Equation

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Received 20, December, 2012

Accepted 12, June, 2013

Abstract:

In this paper we find the exact solution of Burger's equation after reducing it to Bernoulli equation. We compare this solution with that given by Kaya where he used Adomian decomposition method, the solution given by chakrone where he used the Variation iteration method (VIM) and the solution given by Eq(5) in the paper of M. Javidi. We notice that our solution is better than their solutions.

Key words: Burger's equation, kdv equation, PDF.

1. Introduction:

We consider the following Burger's equation

$$u_t + uu_x - \lambda u_{xx} = 0$$

Where λ is positive parameter

This equation arises in various areas of science. Equation (1.1) has been used in the study of the propagation through liquid-filled elastic tube [1] and description for shallow water waves on a viscous fluid [2]. Equation (1.1) is used as a model for traffic flow [3]. Many Researchers had proposed various kinds of exact and numerical solution [4,5,6,7,8,9,10] where they used Adomian decomposition method, variation iteration method, Galerkin method.

2. Derivation of Burger's equation.[9]

We recall the differential form of the nonlinear conservation equation

$$\frac{\partial p}{\partial t} + \frac{\partial q}{\partial x} = 0$$

To investigate the nature of the discontinuous solution or shock waves, we assume a function relation

$q = Q(p)$ and allow a jump discontinuity for p and q .

In many physical problems of interest it would be a better approximation to assume that q is a function of the density gradient p_x as well as p . A

simple model is to take

$$q = Q(p) - vp_x$$

Where v is a positive constant. Substituting (2.2) into (2.1), we obtain the nonlinear diffusion equation

$$p_t + c(p)p_x = vp_{xx}$$

Where,

$$c(p) = Q'(p)$$

We multiply (2.3) by $c'(p)$ to obtain

$$c'(p)p_t + c(p)c'(p)p_x = vc'(p)p_{xx}$$

Hence:

$$c_t + cc_x = vc'(p)p_{xx}$$

And therefore:

$$c_t + cc_x = vc'(p)p_{xx} = vc_{xx} - c''(p)p_x^2 \quad (2.4)$$

If $Q(p)$ is a quadratic function in p , then $c(p)$ is linear in p , and $c''(p) = 0$.

Consequently (2.4) becomes

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$$c_t + cc_x = \nu c_{xx}$$

As a simple model of turbulence, c is replaced by the fluid velocity field $u(x,t)$ to obtain the well-known Burgers equation as

$$u_t + uu_x = \nu u_{xx}$$

Where ν is the kinematic viscosity. Thus, the Burgers equation is a balance between time evaluation, nonlinearity and diffusion.

3. Solution of Burger's equation.

To solve Burger's equation (1.1), we can assume the solution $u(x,t) = v(z)$

Where

$$z = x - \lambda t$$

$-v$ is called the travelling wave solution-Substituting (3.1) into (1.1) becomes

$$v_{zz} = \left(\frac{\nu}{\lambda} - 1\right)v_z$$

Rewrite (3.2) as a system of first order O.D.E's

Let

$$r_1 = v$$

$$r_2 = v_z$$

Where r_1, r_2 represent the velocity and a acceleration respectively.

Then

$$r'_1 = v = r_2$$

$$r'_2 = v_{zz} = \left(\frac{r_1}{\lambda} - 1\right)r_2$$

This system has an infinite number of equilibrium points which is

r_1 - axis

To solve (3.4) divide the two equations to get

$$\frac{dr_2}{dr_1} = \frac{r_1}{\lambda} - 1 \tag{2.5}$$

i.e.

$$dr_2 = \left(\frac{r_1}{\lambda} - 1\right)dr_1$$

$$r_2 = \frac{r_1^2}{2\lambda} - r_1 + c \tag{2.6}$$

Where c is an arbitrary constant.

Hence, from (3.3) we get

$$v_z + v = \frac{v^2}{2\lambda} + c$$

Now, let

$$w = v + k$$

[In order to eliminate c from equation (3.5)] substituting (3.6) into (3.5) becomes

$$w_z + w - k = \frac{(w - k)^2}{2\lambda} + c$$

i.e.

$$w_z + \left(1 + \frac{k}{\lambda}\right)w = \frac{w^2}{2\lambda} + \frac{k^2}{2\lambda} + k + c \tag{3.1}$$

Let $\frac{k^2}{2\lambda} + k + c = 0$ to get

$$\tag{3.2}$$

$$k = -\lambda \mp \sqrt{\lambda^2 - 2\lambda c}$$

Then (3.7) becomes

$$w_z + \left(1 - \frac{k}{\lambda}\right)w = \frac{w^2}{2\lambda} \tag{3.3}$$

Notice that equation (3.9) represents Bernoulli equation. To solve it Suppose

$$\eta = w^{-1} \rightarrow \eta_z = -w^{-2}w_z$$

Substituting (3.10) into (3.9)

$$w^{-2}w_z + \left(1 + \frac{k}{\lambda}\right)w^{-1} = \frac{1}{2\lambda}$$

We get

$$\eta_z - \left(1 + \frac{k}{\lambda}\right)\eta = \frac{-1}{2\lambda}$$

This equation is first-order linear differential equation, its integrating factor is

$$I(z) = e^{-(1+\frac{\kappa}{\lambda})z}$$

And its solution is gives by

$$\eta I = \frac{1}{2(\lambda + \kappa)} e^{-(\frac{\kappa}{\lambda}+1)z} + \alpha$$

Where α is an arbitrary constant so,

$$\eta = (\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})z})$$

From (3.10), we get the solution of Bernoulli equation (3.9)

$$w(z) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})z}}$$

By using (3.6) and (3.11)

Now, from (3.6) $v = w - k$ so the traveling wave solution (3.1) becomes

$$u(x,t) = v(z) = w(z) - k$$

$$u(x,t) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})(x-\lambda t)}} - k$$

Where k is given in (3.8) and α is an arbitrary constant where $k = -\lambda \mp \sqrt{\lambda^2 - 2\lambda c}$ And α, c are arbitrary constants.

4. The solution of Burger's equation with Dirichlet conditions.

From (3.12) the solution of Burger's equation is

$$u(x,t) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})(x-\lambda t)}} - k$$

Dirichlet conditions are

$$BC1:u(0,t) = g(t)$$

$$0 \leq t \leq t_1$$

$$BC2:u(1,t) = h(t)$$

Clearly $u(0, 0) = 0$ (stationary case)

BC1→

$$u(x,t) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{-(1+\frac{\kappa}{\lambda})\lambda t}} - k = g(t)$$

BC2→

$$u(x,t) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})(1-\lambda t)}} - k = h(t)$$

$$u(0,0) = 0 \rightarrow \tag{3.11}$$

$$\frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha} - k = 0$$

Therefore,

$$\alpha = \frac{1}{k} - \frac{1}{2(\lambda + \kappa)} = \frac{2\lambda + k}{2(\lambda + \kappa)k}$$

Remember that from (3.8)

$$k = -\lambda \mp \sqrt{\lambda^2 - \lambda c} \tag{3.12}$$

And from (4.2)

$$h(0) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})}} - k$$

From which we get

$$\alpha = \left[\frac{1}{h(0) + k} - \frac{1}{2(\lambda + \kappa)} \right] e^{-\frac{(\lambda + \kappa)}{\lambda}}$$

From(4.3) and (4.4)we determine α and c

After determining these constants we

find the solution of Burger's equation with Dirichlet conditions.

As a special case let $\lambda = 2$ and $c = 0.75$ then from

(3.8) we get

$$k = -3 \text{ or } -1$$

And from (4.3) we get

$$\alpha = \frac{1}{6} \text{ if } k = -3$$

Or

$$\alpha = \frac{-3}{2} \text{ if } k = -1$$

Case

$$1: \lambda = 2, c = 0.75, k = -3, \alpha = \frac{1}{6}$$

$$u(x,t) = \frac{1}{\frac{-1}{2} + \frac{1}{6}e^{-0.5(x-2t)}} + 3$$

Case2:

$$\lambda = 2, c = 0.75, k = -1, \alpha = \frac{-3}{2}$$

$$u(x,t) = \frac{1}{\frac{1}{2} - \frac{3}{2}e^{0.5(x-2t)}} + 1$$

The graphs of (4.5),(4.6) are in figures (1)and(2) respectively.

The two shapes have the same qualitative behavior but quantitatively different.

5.Comparison.

- The solution of this equation obtained by Kaya

$$u(x,t) = \frac{(\alpha + \beta + (\beta - \alpha)e^{\zeta})}{(1 + e^{\zeta})}$$

Where $\zeta = \frac{\alpha}{\nu}(x - \beta t - \eta)$ where is not easier than our solution.

- The solution obtained by Omar is approximated

solution, while our solution is exact solution.

- The solution obtained by Javidi

$$u(x,t) = \frac{(0.1e^{-A} + 0.5e^{-B} + e^{-C})}{(e^{-A} + e^{-B} + e^{-C})}$$

Where,

$$A = \frac{0.05}{\mu}(x - 0.5 + 4.95t),$$

$$B = \frac{0.05}{\mu}(x - 0.5 + 0.75t) \text{ and}$$

$$C = \frac{0.05}{\mu}(x - 0.5 + 0.375t)$$

Clearly our solution is easier than his solution.

when

$$\alpha = -1.72, \lambda = 0.1, k = 0.2, c = 0 \quad (4.5)$$

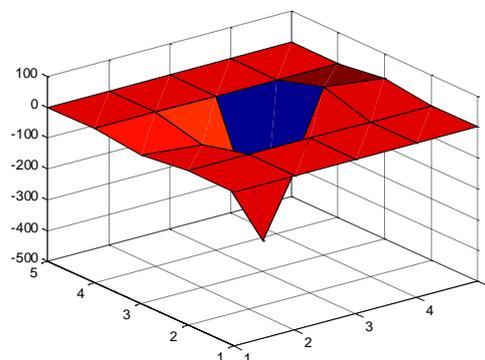
in the solution function

$$u(x,t) = \frac{1}{\frac{1}{2(\lambda + \kappa)} + \alpha e^{(1+\frac{\kappa}{\lambda})(x-\lambda t)}} - k$$

(4.6)

where $t = [0:0.1:0.4], x = [-10:5:10]$

This figure classifies the wave at the beginning of the earth quake.



Case 1: $\lambda = 2, c = 0.75, k = -3, \alpha = \frac{1}{6}$

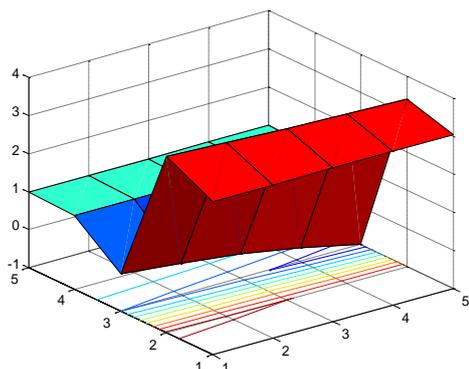


Fig.-1-

Case2: $\lambda = 2, c = 0.75, k = -1, \alpha = \frac{-3}{2}$

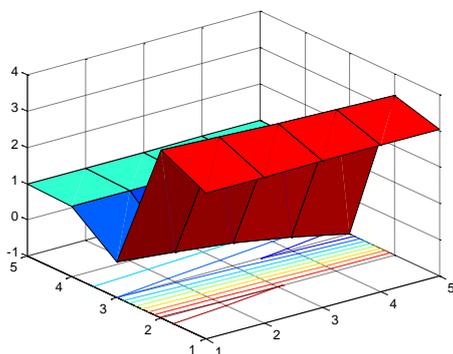


Fig.-2-

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استخدام معادلة برنولي لحل معادلة بيركر

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الخلاصة:

حصلنا من بحثنا هذا على الحل المضبوط (Exact) لمعادلة بيركر بعد تحويلها إلى معادلة برنولي وتمت مقارنة هذا الحل مع حل "Kaya" الذي استخدم طريقة تجزئة ادومين وحل "Chakrone" الذي استخدم طريقة التغيرات المتكرر وحل "Javidi" المعطى بالعلاقة (5) وتبين ان حلنا هذا افضل من حلولهم للبساطه الصيغة